# Parabolic Anderson model with rough noise in space and rough initial conditions* 

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#### Abstract

In this note, we consider the parabolic Anderson model on $\mathbb{R}_{+} \times \mathbb{R}$, driven by a Gaussian noise which is fractional in time with index $H_{0}>1 / 2$ and fractional in space with index $0<H<1 / 2$ such that $H_{0}+H>3 / 4$. Under a general condition on the initial data, we prove the existence and uniqueness of the mild solution and obtain its exponential upper bounds in time for all $p$-th moments with $p \geq 2$.


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## 1 Introduction

In this paper, we study the parabolic Anderson model (PAM):

$$
\left\{\begin{array}{l}
\frac{\partial u}{\partial t}(t, x)=\frac{1}{2} \frac{\partial u}{\partial x^{2}}(t, x)+u(t, x) \dot{W}(t, x) \quad t>0, x \in \mathbb{R},  \tag{1.1}\\
u(0, \cdot)=\mu_{0}
\end{array}\right.
$$

with initial condition given by a non-negative Borel measure $\mu_{0}$ on $\mathbb{R}$ such that

$$
\begin{equation*}
\int_{\mathbb{R}} e^{-a x^{2}} \mu_{0}(\mathrm{~d} x)<\infty \quad \text { for all } a>0 \tag{1.2}
\end{equation*}
$$

Initial conditions of this type (called rough initial conditions) were introduced in [4] and were considered later for the stochastic heat equation in various settings; see, e.g., $[2,5,6]$ and references therein. The noise $\dot{W}$ is assumed to be a centered Gaussian noise that is fractional in time and space with indices $H_{0}$, respectively $H$ in the following range (see the gray area in Figure 1.1):

$$
\begin{equation*}
\left(H_{0}, H\right) \in(1 / 2,1) \times(0,1 / 2) \quad \text { and } \quad H+H_{0}>3 / 4 \tag{1.3}
\end{equation*}
$$

[^0]

Figure 1.1: The gray area - Area II - corresponds to the range of $H_{0}$ and $H$ in which the noise $\dot{W}$ is rough in space. Area I corresponds to the case when the noise is non rough.

Rigorously, $\left\{W(\varphi) ; \varphi \in \mathcal{D}\left(\mathbb{R}_{+} \times \mathbb{R}\right)\right\}$ is a zero-mean Gaussian process with covariance ${ }^{1}$ :

$$
\mathbb{E}[W(\varphi) W(\psi)]=\alpha_{H_{0}} \int_{\mathbb{R}_{+}^{2}} \int_{\mathbb{R}}|t-s|^{2 H_{0}-2} \mathcal{F} \varphi(t, \cdot)(\xi) \overline{\mathcal{F} \psi(s, \cdot)(\xi)} \mu(\mathrm{d} \xi) \mathrm{d} t \mathrm{~d} s=:\langle\varphi, \psi\rangle_{\mathcal{H}}
$$

where $\alpha_{H_{0}}=H_{0}\left(2 H_{0}-1\right)$ and $\mu(\mathrm{d} \xi)=c_{H}|\xi|^{1-2 H} \mathrm{~d} \xi$, with $c_{H}=\Gamma(2 H+1) \sin (\pi H) /(2 \pi)$. Since $H<1 / 2$, we say that $W$ is rough in space.

We denote by $\mathcal{H}$ the completion of $\mathcal{D}\left(\mathbb{R}_{+} \times \mathbb{R}\right)$ with respect to the inner product $\langle\cdot, \cdot\rangle_{\mathcal{H}}$. Then $W=\{W(\varphi)\}_{\varphi \in \mathcal{H}}$ is an isonormal Gaussian process and we can use Malliavin calculus to define and analyze the solution to (1.1). We say that a process $u=\{u(t, x) ; t>$ $0, x \in \mathbb{R}\}$ is a Skorohod solution of (1.1) if it is adapted with respect to the filtration induced by $W$, and for all $t>0$ and $x \in \mathbb{R}$,

$$
\begin{equation*}
u(t, x)=J_{0}(t, x)+\int_{0}^{t} \int_{\mathbb{R}} G(t-s, x-y) u(s, y) W(\delta s, \delta y) \tag{1.4}
\end{equation*}
$$

where $J_{0}$ is the solution to the homogeneous heat equation, i.e.,

$$
\begin{equation*}
J_{0}(t, x):=\int_{\mathbb{R}} G(t, x-y) \mu_{0}(\mathrm{~d} y) \quad \text { with } \quad G(t, x)=(2 \pi t)^{-1 / 2} e^{-\frac{x^{2}}{2 t}} . \tag{1.5}
\end{equation*}
$$

The stochastic integral in (1.4) is interpreted in the Skorohod sense, i.e. it is given by the divergence operator from Malliavin calculus. We refer the reader to Section 1.3 of [12] for the definition of this operator, and to [2, 10, 9] for similar developments.

The following theorem is the main result of the present article.
Theorem 1.1. If $\left(H_{0}, H\right)$ satisfy (1.3) and if $\mu_{0}$ satisfies (1.2), then equation (1.1) has a unique solution $u$ and this solution satisfies: for all $p \geq 2, t>0$ and $x \in \mathbb{R}$,

$$
\begin{equation*}
\mathbb{E}\left(|u(t, x)|^{p}\right) \leq C_{1}^{p} J_{0}^{p}(t, x) \exp \left(C_{2} p^{\frac{H+1}{H}} t^{\frac{2 H_{0}+H-1}{H}}\right) \tag{1.6}
\end{equation*}
$$

where $C_{1}>0$ and $C_{2}>0$ are some constants which depend on $H_{0}$ and $H$.
The novelty of our result is the fact that we consider rough initial condition. One prominent example is the case $\mu_{0}=\delta_{0}$ where $\delta_{0}$ is the Dirac delta measure; see, e.g., [1].
Remark 1.2. When the initial condition is a bounded function or a constant, X. Chen established, in Theorem 1.2 of [7], the well-posedness of the solution to (1.1) in $L^{p}(\Omega)$ under conditions (1.3). Note that recently, Z.-Q. Chen and Y. Hu identified, in part (ii) of

[^1]Theorem 1.3 of [8], the following necessary condition for the well-posedness of (1.1), in the case when $\left(H_{0}, H\right) \in(1 / 2,1) \times(0,1 / 2)$ and $\mu_{0}(\mathrm{~d} x)=u_{0}(x) \mathrm{d} x$ with $u_{0}(x) \geq c>0$ :

$$
H+2 H_{0}>5 / 4 .
$$

In the present paper, we extend X . Chen's work to the rough initial condition case. This extension is highly nontrivial. As we will see in the proof of Theorem 1.1 below, to estimate the moments of the solution, we need to compute a spatial integral and a time integral. In [7], X. Chen uses the Laplace transform to handle the time integral first and then computes the spatial integral. This method does not work for the rough initial data. In a nutshell, moving from bounded initial data to rough initial data, one essentially changes the underlying motion from the Brownian motion to the Brownian bridge. The time increments no longer take the linear difference form; see Lemma 2.1. To overcome this difficulty, we compute the spatial integral first and then use Lemma A. 1 to compute the time integral. This leads to a complicated expression which contains a product $\gamma_{n}$ of ratios of gamma functions; see (2.9). The most delicate part is to estimate $\gamma_{n}$. For this, we develop some novel combinatorial techniques. Unlike [7, 8] where (1.1) in $\mathbb{R}^{d}$ was studied, in this note, we only study on the one-spatial-dimensional case.

Recently, Hu and Lê obtained in Theorem 3.2 of [10] both the well-posedness and the following moment asymptotics: ${ }^{2}$

$$
\begin{equation*}
\mathbb{E}\left(|u(t, x)|^{p}\right) \leq C_{1}^{p} t^{-p\left(\beta+\frac{2 H-1}{4}\right)} \exp \left(C_{2} p^{\frac{H+1}{H}} t^{\frac{2 H_{0}+H-1}{H}}\right), \quad \text { for all } p \geq 2, t>0, x \in \mathbb{R} \tag{1.7}
\end{equation*}
$$

under weaker conditions on $\mu_{0}$, namely, $\mu_{0}$ is a Borel measure such that

$$
\begin{equation*}
\int_{\mathbb{R}}\left(1+|\xi|^{-(H-1 / 2)}\right) e^{-t|\xi|^{2}}\left|\mathcal{F} \mu_{0}(\xi)\right| \mathrm{d} \xi \leq C t^{-\beta} \quad \text { for all } t>0 \tag{1.8}
\end{equation*}
$$

for some $C>0$ and $\beta<H_{0}$, where $\mathcal{F} \mu_{0}$ is the Fourier transform of $\mu_{0}$. Condition (1.8) is more restrictive than (1.2); see Remark 1.3 for one example. While our exponent in (1.6) recovers that of Hu and Lê in (1.7), the factor $J_{0}^{p}(t, x)$ in (1.6) looks more natural than the corresponding factor in (1.7). For example, when the initial condition is a bounded function (resp. the delta initial measure), then as $t \rightarrow 0$, the factor $J_{0}^{p}(t, x)$ will not blow up (resp. blows up at the exact rate $t^{-d / 2}$ for $x=0$ and will not blow up for $x \neq 0$ ). But it is not clear whether the factor $t^{-p(\beta+(2 H-1) / 4)}$ in (1.7) would blow up or not, which depends on the sign of the exponent of $t$.
Remark 1.3. Our condition on the initial data allows growing tails, for example, $\mu(\mathrm{d} x)=$ $x^{2} \mathrm{~d} x$. In this case, $\int_{\mathbb{R}} e^{-a x^{2}} x^{2} \mathrm{~d} x=\sqrt{\pi} 2^{-1} a^{-3 / 2}<\infty$ for all $a>0$. Hence, condition (1.2) is satisfied. But this initial condition cannot satisfy condition (1.8) because $\mathcal{F}\left[x^{2}\right](\xi)=$ $\delta^{\prime \prime}(\xi)$ (in the generalized sense, see, e.g., Theorem 7.4 of [14]), which is a genuine distribution and hence does not have module or absolute value.

Finally, this paper can also be viewed as a continuation of [2] where the case of rough initial conditions and $\left(H_{0}, H\right) \in(1 / 2,1)^{2}$ (see Figure 1.1) was covered.

## 2 Proof of Theorem 1.1

We denote by $I_{n}: \mathcal{H}^{\otimes n} \rightarrow \mathcal{H}$ the multiple Wiener integral of order $n$ with respect to $W$. It is known that the solution $u$ exists if and only if $\sum_{n \geq 1} I_{n}\left(f_{n}(\cdot, t, x)\right)$ converges in

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$L^{2}(\Omega)$, and in this case the solution has the Wiener chaos expansion:

$$
\begin{aligned}
u(t, x) & =J_{0}(t, x)+\sum_{n \geq 1} I_{n}\left(f_{n}(\cdot, t, x)\right) \text { with } \\
f_{n}\left(t_{1}, x_{1}, \ldots, t_{n}, x_{n}, t, x\right) & =\prod_{j=1}^{n} G\left(t_{j+1}-t_{j}, x_{j+1}-x_{j}\right) J_{0}\left(t_{1}, x_{1}\right) 1_{\left\{0<t_{1}<\ldots<t_{n}<t\right\}}
\end{aligned}
$$

and $t_{n+1}=t$ and $x_{n+1}=x$; see for instance [2, 9]. By the orthogonality of the terms in this series, the necessary and sufficient condition for the existence of solution is:

$$
\begin{equation*}
\sum_{n \geq 1} n!\left\|\widetilde{f}_{n}(\cdot, t, x)\right\|_{\mathcal{H}^{\otimes n}}^{2}<\infty \tag{2.1}
\end{equation*}
$$

where $\tilde{f}_{n}(\cdot, t, x)$ is the symmetrization of $f_{n}(\cdot, t, x)$, defined by:

$$
\widetilde{f}_{n}\left(t_{1}, x_{1}, \ldots, t_{n}, x_{n}, t, x\right)=\frac{1}{n!} \sum_{\rho \in S_{n}} f_{n}\left(t_{\rho(1)}, x_{\rho(1)}, \ldots, t_{\rho(n)}, x_{\rho(n)}, t, x\right)
$$

where $S_{n}$ is the set of permutations of $\{1, \ldots, n\}$. For any $\mathbf{t}=\left(t_{1}, \ldots, t_{n}\right) \in[0, t]^{n}$, $\mathbf{s}=\left(s_{1}, \ldots, s_{n}\right) \in[0, t]^{n}$, we denote

$$
\begin{aligned}
\psi_{t, x}^{(n)}(\mathbf{t}, \mathbf{s})=(n!)^{2} \int_{\mathbb{R}^{d}} \mu\left(\mathrm{~d} \xi_{1}\right) \ldots \mu\left(\mathrm{d} \xi_{n}\right) & \mathcal{F} \widetilde{f}_{n}\left(t_{1}, \cdot, \ldots, t_{n}, \cdot, t, x\right)\left(\xi_{1}, \ldots, \xi_{n}\right) \\
& \times \overline{\mathcal{F} \widetilde{f}_{n}\left(s_{1}, \cdot, \ldots, s_{n}, \cdot, t, x\right)\left(\xi_{1}, \ldots, \xi_{n}\right)} .
\end{aligned}
$$

We will use the following result.
Lemma 2.1 (Lemma 3.2 of [2]). If $0<t_{\rho(1)}<\ldots<t_{\rho(n)}<t=: t_{\rho(n+1)}$, then

$$
\psi_{t, x}^{(n)}(\boldsymbol{t}, \boldsymbol{s}) \leq J_{0}^{2}(t, x) \int_{\mathbb{R}^{n}} \prod_{k=1}^{n} \exp \left\{-\frac{t_{\rho(k+1)}-t_{\rho(k)}}{t_{\rho(k+1)} t_{\rho(k)}}\left|\sum_{j=1}^{k} t_{\rho(j)} \xi_{j}\right|^{2}\right\} \mu\left(\mathrm{d} \xi_{1}\right) \ldots \mu\left(\mathrm{d} \xi_{n}\right)
$$

We will also use the following estimate, which is a consequence of Hölder's inequality with $p=1 / H$, and the Littlewood-Hardy-Sobolev inequality (see, e.g., [11]):

$$
\begin{equation*}
\alpha_{H_{0}}^{n} \int_{\mathbb{R}_{+}^{2 n}} \prod_{j=1}^{n}\left|t_{j}-s_{j}\right|^{2 H_{0}-2} \varphi(\mathbf{t}) \varphi(\mathbf{s}) \mathrm{d} \mathbf{t} d \mathbf{s} \leq b_{H_{0}}^{n}\left(\int_{\mathbb{R}_{+}^{n}}|\varphi(\mathbf{t})|^{1 / H_{0}} \mathrm{~d} \mathbf{t}\right)^{2 H_{0}} \tag{2.2}
\end{equation*}
$$

Lemma 2.2. For $n \geq 2$ and $x_{1}, \ldots, x_{n} \in \mathbb{R}_{+}$, it holds that

$$
\begin{equation*}
S_{n}:=x_{1} \prod_{k=2}^{n}\left(x_{k}+x_{k-1}\right)=\sum_{a \in A_{n}} \prod_{j=1}^{n} x_{j}^{a_{j}}, \tag{2.3}
\end{equation*}
$$

where $A_{n}$ is a set of indices $a=\left(a_{1}, \ldots, a_{n}\right)$ such that $\operatorname{card}\left(A_{n}\right)=2^{n-1}$ and

$$
\begin{gather*}
a_{1} \in\{1,2\}, \quad a_{n} \in\{0,1\}, \quad a_{2}, \ldots, a_{n-1} \in\{0,1,2\},  \tag{2.4a}\\
\sum_{j=1}^{i} a_{j} \in\{i, i+1\} \quad \text { for } i=1, \cdots, n-1, \sum_{j=1}^{n} a_{j}=n,  \tag{2.4b}\\
a_{i}+a_{i+1} \in\{1,2,3\} \quad \text { for } i=2, \ldots, n-2,  \tag{2.4c}\\
a_{1}+a_{2} \in\{2,3\} \quad \text { and } \quad a_{n-1}+a_{n} \in\{1,2\} . \tag{2.4d}
\end{gather*}
$$

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Figure 2.2: The eight paths for $n=4$, from $(1,1)$ to either $(4,3)$ or $(4,4)$, correspond to the eight monomials in the expansion of $S_{4}=x_{1}\left(x_{1}+x_{2}\right)\left(x_{2}+x_{3}\right)\left(x_{3}+x_{4}\right)$. All paths should stay in the dashed envelope. The four digits correspond to the value of $\left(a_{1}, \ldots, a_{4}\right)$.

An expansion similar to Lemma 2.2 can be found in [9, p. 488]. Our result is slightly more precise. For example, instead of $a_{1}, a_{n} \in\{0,1,2\}$, we have $a_{1} \in\{1,2\}$ and $a_{n} \in\{0,1\}$.

Proof of Lemma 2.2. Clearly it holds for $n=2$. Assume that the statement is true for $n \geq 2$. Then

$$
S_{n+1}=S_{n}\left(x_{n}+x_{n+1}\right)=\sum_{a \in A_{n}} \prod_{i=1}^{n} x_{i}^{a_{i}} x_{n}^{a_{n}+1}+\sum_{a \in A_{n}} \prod_{i=1}^{n} x_{i}^{a_{i}} x_{n+1}=: \sum_{a^{\prime} \in A_{n+1}} \prod_{i=1}^{n+1} x_{i}^{a_{i}^{\prime}}
$$

The statement for $n+1$ follows by considering separately two cases: (i) $a_{1}^{\prime}=a_{1}, \ldots, a_{n-1}^{\prime}=$ $a_{n-1}, a_{n}^{\prime}=a_{n}+1, a_{n+1}^{\prime}=0$; and (ii) $a_{1}^{\prime}=a_{1}, \ldots,, a_{n}^{\prime}=a_{n}, a_{n+1}^{\prime}=1$.

Remark 2.3. Lemma 2.2 can also be proved using a path representation. Since this representation will be used in the proof of Theorem 1.1, we explain it here. To each monomial $x_{1}^{a_{1}} \ldots x_{n}^{a_{n}}$ in the expansion of $S_{n}$ one can associate a path starting from $(1,1)$ and going to $(n, n)$ or $(n, n-1)$, depending on whether $x_{n}$ is present or absent in the monomial. This path is composed of $n-1$ segments, which correspond to exponents $a_{1}, \ldots, a_{n-1}$ (in this order) and are constructed as follows:
-if the exponent is 0 , the path moves 1 unit to the right and 2 units up;
-if the exponent is 1 , the path moves 1 unit to the right and 1 unit up;
-if the exponent is 2 , the path moves 1 unit to the right and 0 units up.
See Figure 2.2 for an illustration of this correspondence with $n=4$. See also Figure 2.3 for the properties in (2.4).

Now we are ready to prove Theorem 1.1.
Proof of Theorem 1.1. Denote $J_{n}(t, x)=I_{n}\left(f_{n}(\cdot, t, x)\right)$. Let $C$ be a constant which depends on $H_{0}$ and $H$ and may be different from line to line. The proof consists of five steps:

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Figure 2.3: Illustrations of properties in (2.4) with $n=7$. Each monomial in the expansion of $(2.3)$ corresponds to a path from $(1,1)$ to either $(7,7)$ or $(7,6)$. The dots indicate the possible choices for the position of the path.

Step 1. Note that $\psi_{t, x}^{(n)}(\mathbf{t}, \mathbf{s}) \leq \psi_{t, x}^{(n)}(\mathbf{t}, \mathbf{t})^{1 / 2} \psi_{t, x}^{(n)}(\mathbf{s}, \mathbf{s})^{1 / 2}$ by the Cauchy-Schwarz inequality. Combining this with inequality (2.2), we obtain:

$$
\begin{align*}
\mathbb{E}\left(\left|J_{n}(t, x)\right|^{2}\right) & =n!\left\|\tilde{f}_{n}(\cdot, t, x)\right\|_{\mathcal{H} \otimes n}^{2}=\frac{1}{n!} \alpha_{H_{0}}^{n} \int_{[0, t]^{2 n}} \prod_{j=1}^{n}\left|t_{j}-s_{j}\right|^{2 H_{0}-2} \psi_{t, x}^{(n)}(\mathbf{t}, \mathbf{s}) \mathrm{d} \mathbf{t} \mathrm{~d} \mathbf{s} \\
& \leq \frac{1}{n!} b_{H_{0}}^{n}\left(\sum_{\rho \in S_{n}} \int_{0<t_{\rho(1)}<\ldots<t_{\rho(n)}<t} \psi_{t, x}^{(n)}(\mathbf{t}, \mathbf{t})^{\frac{1}{2 H_{0}}} \mathrm{~d} \mathbf{t}\right)^{2 H_{0}} \tag{2.5}
\end{align*}
$$

By Lemma 2.1, for any $\rho \in S_{n}$ fixed, we have:

$$
\begin{aligned}
& \int_{0<t_{\rho(1)}<\ldots<t_{\rho(n)}<t} \psi_{t, x}^{(n)}(\mathbf{t}, \mathbf{t})^{\frac{1}{2 H_{0}}} \mathrm{~d} \mathbf{t} \leq J_{0}^{1 / H_{0}}(t, x) \int_{0<t_{\rho(1)}<\ldots<t_{\rho(n)}<t} \mathrm{~d} \mathbf{t} \\
& \quad \times\left(\int_{\mathbb{R}^{n}} \prod_{k=1}^{n} \exp \left\{-\frac{t_{\rho(k+1)}-t_{\rho(k)}}{t_{\rho(k+1)} t_{\rho(k)}}\left|\sum_{j=1}^{k} t_{\rho(j)} \xi_{j}\right|^{2}\right\} \mu\left(\mathrm{d} \xi_{1}\right) \ldots \mu\left(\mathrm{d} \xi_{n}\right)\right)^{\frac{1}{2 H_{0}}} .
\end{aligned}
$$

Using the change of variables $t_{k}^{\prime}=t_{\rho(k)}$ for $k=1, \ldots, n$, we see that the integral on the right-hand side above does not depend on $\rho$. Hence,

$$
\int_{0<t_{\rho(1)}<\ldots<t_{\rho(n)}<t} \psi_{t, x}^{(n)}(\mathbf{t}, \mathbf{t})^{\frac{1}{2 H_{0}}} \mathrm{~d} \mathbf{t} \leq J_{0}^{1 / H_{0}}(t, x) \int_{\left\{0<t_{1}<\ldots<t_{n}<t\right\}} I_{t}^{(n)}\left(t_{1}, \ldots, t_{n}\right)^{\frac{1}{2 H_{0}}} \mathrm{~d} \mathbf{t},
$$

where

$$
I_{t}^{(n)}\left(t_{1}, \ldots, t_{n}\right)=\int_{\mathbb{R}^{n}} \prod_{k=1}^{n} \exp \left\{-\frac{t_{k+1}-t_{k}}{t_{k+1} t_{k}}\left|\sum_{j=1}^{k} t_{j} \xi_{j}\right|^{2}\right\} \mu\left(\mathrm{d} \xi_{1}\right) \ldots \mu\left(\mathrm{d} \xi_{n}\right)
$$

and $t_{n+1}=t$. Taking the sum over all $\rho \in S_{n}$ and coming back to (2.5), we obtain:

$$
\begin{equation*}
\mathbb{E}\left(\left|J_{n}(t, x)\right|^{2}\right) \leq J_{0}^{2}(t, x) b_{H_{0}}^{n}(n!)^{2 H_{0}-1}\left(\int_{\left\{0<t_{1}<\ldots<t_{n}<t\right\}} I_{t}^{(n)}\left(t_{1}, \ldots, t_{n}\right)^{\frac{1}{2 H_{0}}} \mathrm{~d} \mathbf{t}\right)^{2 H_{0}} \tag{2.6}
\end{equation*}
$$

This inequality is similar to Lemma 3.3 of [2].
Step 2. We now estimate $I_{t}^{(n)}\left(t_{1}, \ldots, t_{n}\right)$. We use the change of variables $z_{j}=t_{j} \xi_{j}$ for $j=1, \ldots, n$, followed by $\eta_{k}=\sum_{j=1}^{k} z_{j}$ for $k=1, \ldots, n$. We obtain:

$$
\begin{aligned}
& I_{t}^{(n)}\left(t_{1}, \ldots, t_{n}\right)=c_{H}^{n}\left(\prod_{i=1}^{n} t_{i}\right)^{2 H-2} \int_{\mathbb{R}^{n}} \mathrm{~d} \vec{z} \prod_{k=1}^{n} \exp \left\{-\frac{t_{k+1}-t_{k}}{t_{k+1} t_{k}}\left|\sum_{j=1}^{k} z_{j}\right|^{2}\right\}\left|z_{k}\right|^{1-2 H} \\
& =c_{H}^{n}\left(\prod_{i=1}^{n} t_{i}\right)^{2 H-2} \int_{\mathbb{R}^{n}} \mathrm{~d} \vec{\eta}\left(\prod_{k=1}^{n} \exp \left\{-\frac{t_{k+1}-t_{k}}{t_{k+1} t_{k}}\left|\eta_{k}\right|^{2}\right\}\right)\left|\eta_{1}\right|^{1-2 H} \prod_{k=2}^{n}\left|\eta_{k}-\eta_{k-1}\right|^{1-2 H} \\
& \leq c_{H}^{n}\left(\prod_{i=1}^{n} t_{i}\right)^{2 H-2} \int_{\mathbb{R}^{n}} \mathrm{~d} \vec{\eta}\left(\prod_{k=1}^{n} \exp \left\{-\frac{t_{k+1}-t_{k}}{t_{k+1} t_{k}}\left|\eta_{k}\right|^{2}\right\}\right)\left|\eta_{1}\right|^{1-2 H} \prod_{k=2}^{n}\left(\left|\eta_{k}\right|^{1-2 H}\left|+\left|\eta_{k-1}\right|^{1-2 H}\right),\right.
\end{aligned}
$$

where $\mathrm{d} \vec{z}=\mathrm{d} z_{1} \ldots \mathrm{~d} z_{n}$ and similarly $\mathrm{d} \vec{\eta}=\mathrm{d} \eta_{1} \ldots \mathrm{~d} \eta_{n}$.
By Lemma 2.2 (see also Remark 2.3 and Figure 2.2 for more explanations),

$$
\left|\eta_{1}\right|^{1-2 H} \prod_{k=2}^{n}\left(\left|\eta_{k}\right|^{1-2 H}\left|+\left|\eta_{k-1}\right|^{1-2 H}\right)=\sum_{a \in A_{n}} \prod_{j=1}^{n}\left|\eta_{j}\right|^{(1-2 H) a_{j}}=\sum_{\alpha \in D_{n}} \prod_{j=1}^{n}\left|\eta_{j}\right|^{\alpha_{j}},\right.
$$

where $D_{n}$ is the set of all multi-indices $\alpha=\left(\alpha_{1}, \ldots, \alpha_{n}\right)$ with $\alpha_{j}=(1-2 H) a_{j}$ for all $j=1, \ldots, n$, and $a=\left(a_{1}, \ldots, a_{n}\right) \in A_{n}$. Therefore,

$$
I_{t}^{(n)}\left(t_{1}, \ldots, t_{n}\right) \leq c_{H}^{n}\left(\prod_{j=1}^{n} t_{j}\right)^{2 H-2} \sum_{\alpha \in D_{n}} \prod_{j=1}^{n}\left\{\int_{\mathbb{R}} \exp \left(-\frac{t_{j+1}-t_{j}}{t_{j+1} t_{j}}\left|\eta_{j}\right|^{2}\right)\left|\eta_{j}\right|^{\alpha_{j}} d \eta_{j}\right\}
$$

Each of the integrals above can be computed explicitly. By Lemma 3.1 of [3],

$$
\int_{\mathbb{R}} e^{-t|\xi|^{2}}|\xi|^{\alpha} \mathrm{d} \xi=\Gamma\left(\frac{1+\alpha}{2}\right) t^{-\frac{1+\alpha}{2}} \quad \text { for any } t>0 \text { and } \alpha>-1
$$

Hence,

$$
\begin{aligned}
I_{t}^{(n)}\left(t_{1}, \ldots, t_{n}\right) & \leq C^{n}\left(\prod_{j=1}^{n} t_{j}\right)^{2 H-2} \sum_{\alpha \in D_{n}} \prod_{j=1}^{n}\left(\frac{t_{j+1}-t_{j}}{t_{j} t_{j+1}}\right)^{-\frac{1+\alpha_{j}}{2}} \\
& =C^{n} \sum_{\alpha \in D_{n}} t_{1}^{\frac{4 H-3+\alpha_{1}}{2}}\left(\prod_{j=2}^{n} t_{j}^{\frac{4 H-2+\alpha_{j-1}+\alpha_{j}}{2}}\right) t^{\frac{\alpha_{n}+1}{2}} \prod_{j=1}^{n}\left(t_{j+1}-t_{j}\right)^{-\frac{\alpha_{j}+1}{2}} .
\end{aligned}
$$

Step 3. Taking power $1 /\left(2 H_{0}\right)$ and returning to (2.6), we obtain:

$$
\begin{equation*}
J_{n}(t, x) \leq J_{0}^{2}(t, x) C^{n}(n!)^{2 H_{0}-1}\left(\sum_{\alpha \in D_{n}} t^{\frac{\alpha_{n}+1}{4 H_{0}}} \int_{\left\{0<t_{1}<\ldots<t_{n}<t\right\}} \prod_{i=1}^{n} t_{i}^{\widetilde{\alpha}_{i}}\left(t_{i+1}-t_{i}\right)^{\widetilde{\beta}_{i}} \mathrm{~d} \mathbf{t}\right)^{2 H_{0}} \tag{2.7}
\end{equation*}
$$

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where

$$
\widetilde{\alpha}_{j}=\left\{\begin{array}{ll}
\frac{4 H-3+\alpha_{1}}{4 H_{0}} & j=1, \\
\frac{4 H-2+\alpha_{j-1}+\alpha_{j}}{4 H_{0}} & j=2, \ldots, n,
\end{array} \quad \text { and } \quad \widetilde{\beta}_{j}=-\frac{\alpha_{j}+1}{4 H_{0}}, \quad j=1, \ldots, n .\right.
$$

Now we verify that the conditions of Lemma A. 1 hold for the integrals in (2.7). Clearly, $\widetilde{\alpha}_{1}>-1$. When $\alpha_{j}=2(1-2 H)$, condition $\widetilde{\beta}_{j}>-1$ becomes $H+H_{0}>3 / 4$; see (1.3). Now we verify that

$$
\begin{equation*}
\sum_{i=1}^{k}\left(\widetilde{\alpha}_{i}+\widetilde{\beta}_{i}\right)+k+1+\alpha_{k+1}>0 \quad \text { for all } k=1, \ldots, n-1, \tag{2.8}
\end{equation*}
$$

using induction on $k$. For $k=1$, using again the condition $H_{0}+H>3 / 4$, we see that

$$
\widetilde{\alpha}_{1}+\widetilde{\beta}_{1}+\widetilde{\alpha}_{2}+2=\frac{8 H_{0}+8 H-6+\alpha_{1}+\alpha_{2}}{4 H_{0}}>0
$$

Suppose that (2.8) holds for $k-1$. We write

$$
\sum_{i=1}^{k}\left(\widetilde{\alpha}_{i}+\widetilde{\beta}_{i}\right)+k+1+\alpha_{k+1}=\left(\sum_{i=1}^{k-1}\left(\widetilde{\alpha}_{i}+\widetilde{\beta}_{i}\right)+\widetilde{\alpha}_{k}+k\right)+\left(\alpha_{k+1}+\widetilde{\beta}_{k}+1\right)
$$

and we notice that $\alpha_{k+1}+\widetilde{\beta}_{k}+1=\left(4 H_{0}+4 H-3+\alpha_{k+1}\right) /\left(4 H_{0}\right)>0$. Therefore, we can apply Lemma A. 1 to see that

$$
\begin{align*}
& \int_{\left\{0<t_{1}<\ldots<t_{n}<t\right\}} \prod_{i=1}^{n} t_{i}^{\widetilde{\alpha}_{i}}\left(t_{i+1}-t_{i}\right)^{\widetilde{\beta}_{i}} \mathrm{~d} \mathbf{t}=\frac{\Gamma\left(\widetilde{\alpha}_{1}+1\right) \prod_{i=1}^{n} \Gamma\left(\widetilde{\beta}_{i}+1\right)}{\Gamma(|\widetilde{\alpha}|+|\widetilde{\beta}|+n+1)} \gamma_{n} t^{|\widetilde{\alpha}|+|\widetilde{\beta}|+n} \\
& \quad \text { with } \quad \gamma_{n}:=\prod_{k=1}^{n-1} \frac{\Gamma\left(\sum_{i=1}^{k}\left(\widetilde{\alpha}_{i}+\widetilde{\beta}_{i}\right)+k+1+\widetilde{\alpha}_{k+1}\right)}{\Gamma\left(\sum_{i=1}^{k}\left(\widetilde{\alpha}_{i}+\widetilde{\beta}_{i}\right)+k+1\right)} \tag{2.9}
\end{align*}
$$

Step 4. In this step, we will show that $\gamma_{n} \leq 1$. Note that

$$
\widetilde{\alpha}_{1}+\widetilde{\beta}_{1}=\frac{H-1}{H_{0}} \quad \text { and } \quad \widetilde{\alpha}_{k}+\widetilde{\beta}_{k}=\frac{4 H-3+\alpha_{k-1}}{4 H_{0}}, \quad k=2, \ldots, n
$$

Denote $\theta_{k}:=\sum_{i=1}^{k}\left(\widetilde{\alpha}_{i}+\widetilde{\beta}_{i}\right)+k+1$. Hence,

$$
\theta_{k}= \begin{cases}\frac{H-1}{H_{0}}+2 & \text { if } k=1  \tag{2.10}\\ 1-\frac{1}{4 H_{0}}+k \frac{4 H_{0}+4 H-3}{4 H_{0}}+\frac{1-2 H}{4 H_{0}} \sum_{i=1}^{k-1} a_{i} & \text { if } k=2, \ldots, n\end{cases}
$$

Note that

$$
\begin{equation*}
\theta_{k}-\theta_{k-1}=\frac{4 H_{0}+4 H-3}{4 H_{0}}+\frac{1-2 H}{4 H_{0}} a_{k-1}, \quad \text { for } k=2, \cdots, n \tag{2.11}
\end{equation*}
$$

We see that $\gamma_{n}$ is a function of $a_{i}$ :

$$
\gamma_{n}\left(a_{1}, \ldots, a_{n}\right)=\prod_{k=1}^{n-1} \frac{\Gamma\left(\theta_{k}+\frac{1-2 H}{4 H_{0}}\left(a_{k}+a_{k+1}-2\right)\right)}{\Gamma\left(\theta_{k}\right)}
$$

Recall that any choice of $a_{i}$ corresponds to a path as shown in Figure 2.3. We claim that when we move the path downwards, the value of $\gamma_{n}$ decreases.

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| $k$ | $\theta_{k}^{\prime}-\theta_{k}$ | $\left(a_{k}^{\prime}+a_{k+1}^{\prime}\right)-\left(a_{k}+a_{k+1}\right)$ | $a_{k}^{\prime}-a_{k}$ | $\stackrel{1}{1}$ | Case I |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $i+2$ | 0 | 0 | 0 | $i+2-$ | $x_{i+2}$ |
| $i+1$ | $\frac{1-2 H}{4 H_{0}}$ | -1 | -1 | $i+1-$ |  |
| $i$ | 0 | 0 | 1 | $i-$ | $x_{i}$--- ${ }^{\circ} x_{i}$ |
| $i-1$ | 0 | 1 | 0 | $i-1$ - | $x_{i-1}^{\bullet}$ |





Figure 2.4: Four cases for the positions of the paths $\left(a_{1}, \ldots, a_{n}\right)$ and $\left(a_{1}^{\prime}, \ldots, a_{n}^{\prime}\right)$ around point $(i+1, i+1)$. The values of $a_{k}^{\prime}-a_{k},\left(a_{k}^{\prime}+a_{k+1}^{\prime}\right)-\left(a_{k}+a_{k+1}\right)$, and $\theta_{k}^{\prime}-\theta_{k}$ are the same for all four cases.

As a consequence, the path that achieves the maximum for $\gamma_{n}$ is the one going through $(i, i)$ for $i=1, \ldots, n$, i.e., the straight diagonal line - the topmost line. In this case, we have all $a_{i}$ are equal to one and hence $a_{i}+a_{i+1}=2$ for all $i=1, \ldots, n-1$. Therefore,

$$
\begin{equation*}
\gamma_{n}\left(a_{1}, \ldots, a_{n}\right) \leq \gamma(1, \ldots, 1)=1 \tag{2.13}
\end{equation*}
$$

It remains to prove the claim (2.12). Note that all paths stay between the diagonal and the line parallel to the diagonal, one unit down. If the path does not touch the diagonal, then no action is taken (the argument below will show that the value of $\gamma_{n}$ is minimal for this path).

Let $\left(a_{1}, \ldots, a_{n}\right) \in A_{n}$ be a path which touches the diagonal on at least one point. Say this point is $(i+1, i+1)$ with $i+1<n$. (The case $i+1=n$ is similar.) We compare the value $\gamma_{n}\left(a_{1}, \ldots, a_{n}\right)$ with the value $\gamma_{n}\left(a_{1}^{\prime}, \ldots, a_{n}^{\prime}\right)$ corresponding to another path $\left(a_{1}^{\prime}, \ldots, a_{n}^{\prime}\right) \in A_{n}$ which is obtained by moving the point $(i+1, i+1) 1$ unit down. There are 4 possible cases for the shapes of the two paths around the point $(i+1, i+1)$, which are illustrated in Figure 2.4. Since $a_{k}$ gives the number of points that the path $\left(a_{1}, \ldots, a_{n}\right)$ has on line $k$, it follows that in all 4 cases, $a_{i}^{\prime}=a_{i}+1$ (since line $i$ received one point), $a_{i+1}^{\prime}=a_{i+1}-1$ (since line $i+1$ lost one point) and $a_{k}^{\prime}=a_{k}$ for all $k \notin\{i, i+1\}$ (since the rest of the path remains unchanged). Hence,

$$
\left(a_{k}^{\prime}+a_{k+1}^{\prime}\right)-\left(a_{k}+a_{k+1}\right)=\left\{\begin{array}{ll}
1 & \text { if } k=i-1 \\
-1 & \text { if } k=i+1 \\
0 & \text { otherwise }
\end{array} \quad \text { and } \quad \theta_{k}^{\prime}-\theta_{k}= \begin{cases}\frac{1-2 H}{4 H_{0}} & \text { if } k=i+1 \\
0 & \text { otherwise }\end{cases}\right.
$$

By direct calculation, we see that

$$
\begin{aligned}
\theta_{i-1}^{\prime}+\frac{1-2 H}{4 H_{0}}\left(a_{i-1}^{\prime}+a_{i}^{\prime}-2\right) & =\theta_{i-1}+\frac{1-2 H}{4 H_{0}}\left(a_{i-1}+a_{i}-2\right)+\frac{1-2 H}{4 H_{0}} \text { and } \\
\theta_{k}^{\prime}+\frac{1-2 H}{4 H_{0}}\left(a_{k}^{\prime}+a_{k+1}^{\prime}-2\right) & =\theta_{k}+\frac{1-2 H}{4 H_{0}}\left(a_{k}+a_{k+1}-2\right) \quad \text { for all } k \neq i-1 .
\end{aligned}
$$

Therefore,

$$
\begin{aligned}
\frac{\gamma\left(a_{1}^{\prime}, \cdots, a_{n}^{\prime}\right)}{\gamma\left(a_{1}, \cdots, a_{n}\right)} & =\prod_{k=1}^{n-1} \frac{\Gamma\left(\theta_{k}^{\prime}+\frac{1-2 H}{4 H_{0}}\left(a_{k}^{\prime}+a_{k+1}^{\prime}-2\right)\right)}{\Gamma\left(\theta_{k}+\frac{1-2 H}{4 H_{0}}\left(a_{k}+a_{k+1}-2\right)\right)} \times \prod_{k=1}^{n-1} \frac{\Gamma\left(\theta_{k}\right)}{\Gamma\left(\theta_{k}^{\prime}\right)} \\
& =\frac{\Gamma\left(\theta_{i-1}+\frac{1-2 H}{4 H_{0}}\left(a_{i-1}+a_{i}-2\right)+\frac{1-2 H}{4 H_{0}}\right)}{\Gamma\left(\theta_{i-1}+\frac{1-2 H}{4 H_{0}}\left(a_{i-1}+a_{i}-2\right)\right)} \times \frac{\Gamma\left(\theta_{i+1}\right)}{\Gamma\left(\theta_{i+1}+\frac{1-2 H}{4 H_{0}}\right)} .
\end{aligned}
$$

By applying Lemma A.2, we see that the above ratio is always less than or equal to one. For this, we need to check that

$$
z_{1}:=\theta_{i-1}+\frac{1-2 H}{4 H_{0}}\left(a_{i-1}+a_{i}-2\right) \leq z_{2}:=\theta_{i+1} .
$$

This is clear, since by (2.11), $\theta_{i+1}=\theta_{i-1}+2 \frac{4 H_{0}+4 H-3}{4 H_{0}}+\frac{1-2 H}{4 H_{0}}\left(a_{i-1}+a_{i}-2\right)$. Here we use again condition (1.3). This proves the claim in (2.12).

Step 5. We claim that for $n$ large enough,

$$
\begin{equation*}
\Gamma(|\widetilde{\alpha}|+|\widetilde{\beta}|+n+1) \geq C^{n}(n!)^{\frac{2 H_{0}+H-1}{2 H_{0}}} \tag{2.14}
\end{equation*}
$$

Indeed, by (2.10),

$$
|\widetilde{\alpha}|+|\widetilde{\beta}|+n+1=n \frac{2 H_{0}+H-1}{2 H_{0}}-\frac{1+\alpha_{n}}{4 H_{0}}+1 \geq n \frac{2 H_{0}+H-1}{2 H_{0}}-\frac{1-H}{2 H_{0}}+1 .
$$

We use the fact that for any $a>0, b \in \mathbb{R}$, there exists $N_{a, b} \in \mathbb{N}$ depending on $a$ and $b$ such that $\Gamma(a n+1+b) \geq C_{a, b}^{n}(n!)^{a}$ for all $n \geq N_{a, b}$. Since $\Gamma$ is increasing on $(2, \infty)$, we see that (2.14) holds true. Therefore, thanks to (2.14), we see that for $n$ large enough,

$$
\int_{\left\{0<t_{1}<\ldots<t_{n}<t\right\}} \prod_{i=1}^{n} t_{i}^{\widetilde{\alpha}_{i}}\left(t_{i+1}-t_{i}\right)^{\widetilde{\beta}_{i}} \mathrm{~d} \mathbf{t} \leq C^{n}(n!)^{-\frac{2 H_{0}+H-1}{2 H_{0}}} t^{\frac{2 H_{0}+H-1}{2 H_{0}}-\frac{1+\alpha_{n}}{4 H_{0}}} .
$$

Returning to (2.7), it follows that $\mathbb{E}\left(\left|J_{n}(t, x)\right|^{2}\right) \leq J_{0}^{2}(t, x) C^{n}(n!)^{-H} t^{n\left(2 H_{0}+H-1\right)}$. Finally, by hypercontractivity, the $\|\cdot\|_{p}$-norm on $L^{p}(\Omega)$ is equivalent to the $\|\cdot\|_{2}$-norm (see e.g. page 62 of [12]), and hence

$$
\begin{aligned}
\|u(t, x)\|_{p} & \leq \sum_{n \geq 0}(p-1)^{n / 2}\left\|J_{n}(t, x)\right\|_{2} \leq J_{0}(t, x) \sum_{n \geq 0}(p-1)^{n / 2} C^{n / 2} \frac{1}{(n!)^{H / 2}} t^{n \frac{2 H_{0}+H-1}{2}} \\
& \leq C \exp \left(C p^{1 / H} t^{\frac{2 H_{0}+H-1}{H}}\right),
\end{aligned}
$$

where for the last line we used the fact that $\sum_{n \geq 0} \frac{x^{n}}{(n!)^{a}} \leq C \exp \left(c x^{1 / a}\right)$ for any $x>0$ and $a>0$. We conclude the proof of Theorem 1.1 by taking power $p$.

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## A Some auxiliary results

Lemma A.1. Suppose that $\alpha_{1}>-1, \beta_{i}>-1$ for any $i=1, \ldots, n$, and

$$
\begin{equation*}
\sum_{i=1}^{k}\left(\alpha_{i}+\beta_{i}\right)+k+1+\alpha_{k+1}>0 \quad \text { for all } k=1, \ldots, n-1 \tag{A.1}
\end{equation*}
$$

Then by setting $t_{n+1}=t,|\alpha|=\sum_{i=1}^{n} \alpha_{i}$ and $|\beta|=\sum_{i=1}^{n} \beta_{i}$, we have that

$$
\begin{align*}
& I_{n}\left(t, \alpha_{1}, \ldots, \alpha_{n}, \beta_{1}, \ldots, \beta_{n}\right):=\int_{\left\{0<t_{1}<\ldots<t_{n}<t\right\}} \prod_{i=1}^{n} t_{i}^{\alpha_{i}}\left(t_{i+1}-t_{i}\right)^{\beta_{i}} \mathrm{~d} \boldsymbol{t} \\
& \quad=\frac{\Gamma\left(\alpha_{1}+1\right) \prod_{i=1}^{n} \Gamma\left(\beta_{i}+1\right)}{\Gamma(|\alpha|+|\beta|+n+1)} \prod_{k=1}^{n-1} \frac{\Gamma\left(\sum_{i=1}^{k}\left(\alpha_{i}+\beta_{i}\right)+k+1+\alpha_{k+1}\right)}{\Gamma\left(\sum_{i=1}^{k}\left(\alpha_{i}+\beta_{i}\right)+k+1\right)} t^{|\alpha|+|\beta|+n} . \tag{A.2}
\end{align*}
$$

Proof. The lemma is proved by induction. For $n=1$, we have:

$$
I_{1}(t, \alpha, \beta)=\int_{0}^{t} t_{1}^{\alpha_{1}}\left(t-t_{1}\right)^{\beta_{1}} \mathrm{~d} t_{1}=\frac{\Gamma\left(\alpha_{1}+1\right) \Gamma\left(\beta_{1}+1\right)}{\Gamma\left(\alpha_{1}+\beta_{1}+2\right)} t^{\alpha_{1}+\beta_{1}+1}
$$

For the induction step, we use the fact that

$$
I_{n}\left(t, \alpha_{1}, \ldots, \alpha_{n}, \beta_{1}, \ldots, \beta_{n}\right)=\int_{0}^{t} t_{n}^{\alpha_{n}}\left(t-t_{n}\right)^{\beta_{n}} I_{n-1}\left(t_{n}, \alpha_{1}, \ldots, \alpha_{n-1}, \beta_{1}, \ldots, \beta_{n-1}\right) \mathrm{d} t_{n}
$$

This proves the lemma.
Lemma A.2. For any $a>0$, the function $z \mapsto \Gamma(z+a) / \Gamma(z)$ is non-decreasing on $(0, \infty)$.
Proof. Let $f(z)=\Gamma(z+a) / \Gamma(z)$. Note that $f^{\prime}(z)=f(z)(\psi(z+a)-\psi(z))$, where $\psi(z):=$ $\Gamma^{\prime}(z) / \Gamma(z)$ is the psi function; see 5.2.2 on p. 136 of [13]. By the following expression

$$
\psi(z)=-\gamma+\sum_{k \geq 1}\left(\frac{1}{k+1}-\frac{1}{k+z}\right) \quad \text { for any } \quad z>0
$$

with $\gamma=\lim _{n \rightarrow \infty}\left(\sum_{k=1}^{n} k^{-1}-\ln n\right) \approx 0.5772$ being the Euler's constant (see, e.g., 5.7.6 on p. 139 ibid.), we see that $\psi$ is a nondecreasing function on $(0, \infty)$. Hence, $f^{\prime}(z) \geq 0$ for all $z>0$, which implies the desired result.

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[^1]:    ${ }^{1}$ In this note, we denote by $\mathcal{F} \varphi=\int_{\mathbb{R}^{d}} e^{-i \xi \cdot x} \varphi(x) \mathrm{d} x$ the Fourier transform of a function $\varphi \in L^{1}\left(\mathbb{R}^{d}\right)$.

[^2]:    ${ }^{2}$ This is relation (3.6) of [10] with $\alpha_{0}=2-2 H_{0}$ and $\alpha=2 H-1$; see also Remark 3.5 (ii) ibid.

